

# Dinamika

# Dinamika krutog tela

Mehanika  
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# Literatura

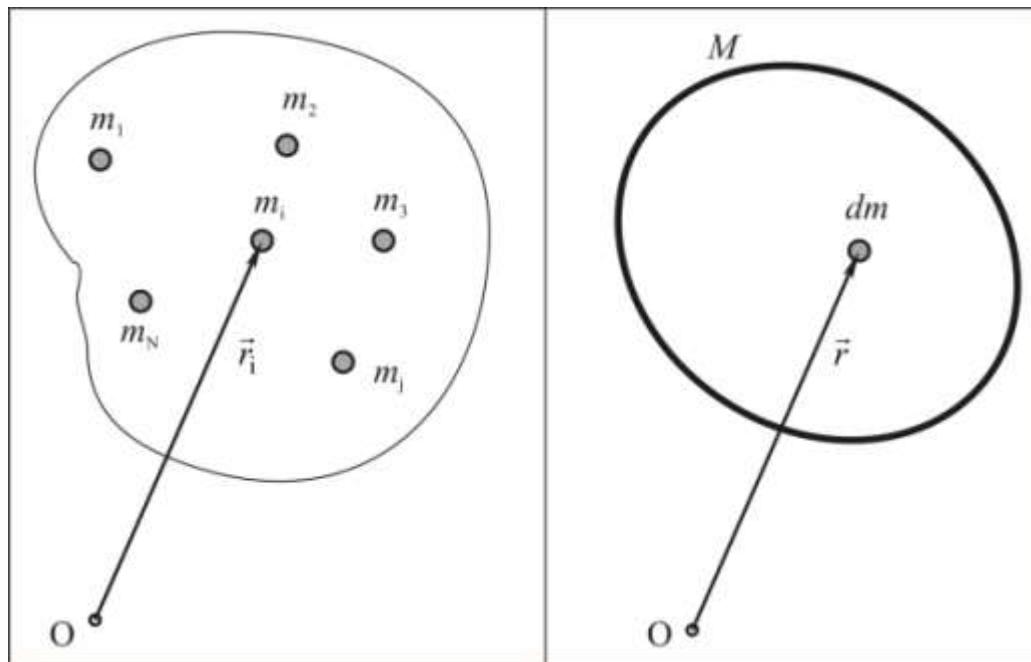
- Đorđe S. Đukić, Teodor M. Atanacković, Ljilja J. Cvetićanin: **Mehanika**, Fakultet tehničkih nauka u Novom Sadu, Novi Sad, 2003.



# Šta ćemo naučiti?

- Translatorno kretanje
- Obrtanje oko nepokretne ose
- Ravansko kretanje

# Sistem tačaka – kruto telo



$$a = \sum_{i=1}^N m_i b_i \rightarrow a = \int_V b \, dm$$

$$M = \sum_{i=1}^N m_i \rightarrow M = \int_V dm$$

$$M\vec{r}_C = \sum_{i=1}^N m_i \vec{r}_i \rightarrow M\vec{r}_C = \int_V \vec{r} \, dm$$

$$\vec{L}_O = \sum_{i=1}^N \vec{L}_{O_i} = \sum_{i=1}^N \vec{r}_i \times m_i \vec{v}_i \rightarrow$$

$$\vec{L}_O = \int_V d\vec{L}_O = \int_V \vec{r} \times \vec{v} \, dm$$

$$E_K = \sum_{i=1}^N E_{Ki} = \sum_{i=1}^N \frac{1}{2} m_i v_i^2 \rightarrow$$

$$E_K = \int_V \frac{1}{2} v^2 dm$$

# Translatorno kretanje

- Diferencijalna jednačina kretanja

$$M\vec{a}_C = \vec{F}_g^s$$

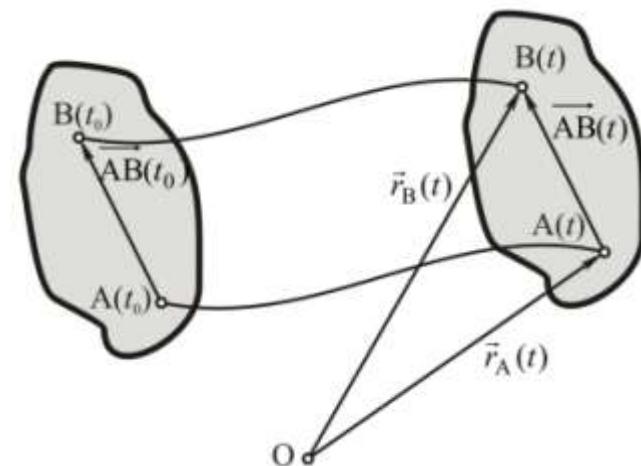
- Zakon o promeni momenta količine kretanja (uslov translatornog kretanja)

$$\vec{v}_r = \vec{v}_i - \vec{v}_C = 0$$

$$d\vec{L}_C = \vec{\rho} \times dm \vec{v}_r = 0$$

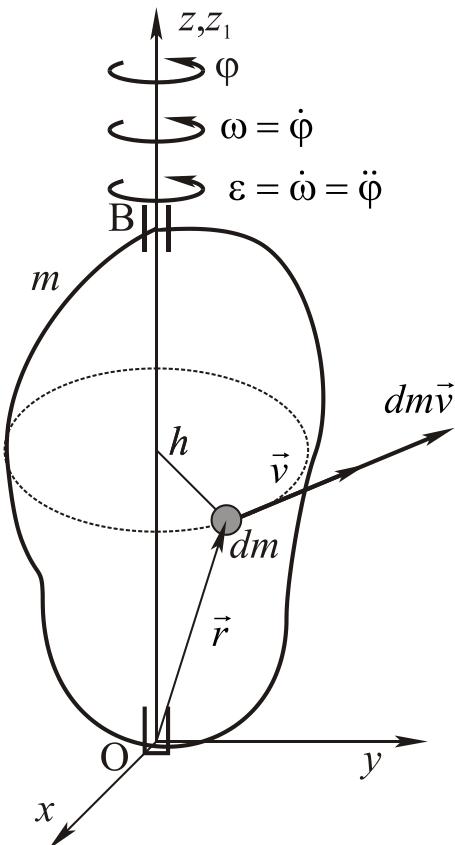
$$\vec{L}_C \equiv 0 \rightarrow \dot{\vec{L}}_C \equiv 0$$

$$\dot{\vec{L}}_C = \overline{\mathbf{M}}_{gC}^s \rightarrow \overline{\mathbf{M}}_{gC}^s = 0$$



# Obrtanje oko nepokretne ose

- Diferencijalna jednačina kretanja



$$\frac{d\vec{L}_O}{dt} = \vec{M}_O^s \rightarrow \boxed{\frac{dL_z}{dt} = \sum M_z^s}$$

$$dL_z = hdmv = h^2 \omega dm \rightarrow L_z = \omega \int_B h^2 dm$$

$$\boxed{J_z = \int_B h^2 dm}$$

$$\boxed{L_z = J_z \omega = J_z \dot{\phi}}$$

$$\boxed{\frac{dL_z}{dt} = J_z \dot{\omega} = J_z \dot{\phi} = J_z \ddot{\phi}}$$

$$J_z \varepsilon = \sum M_z^s$$

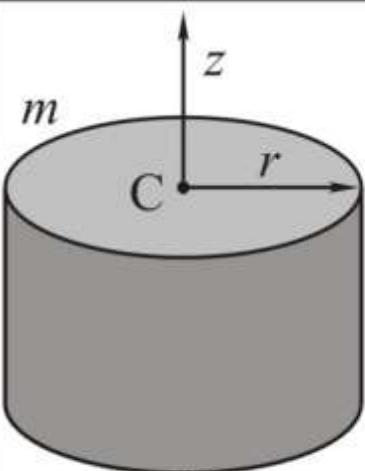
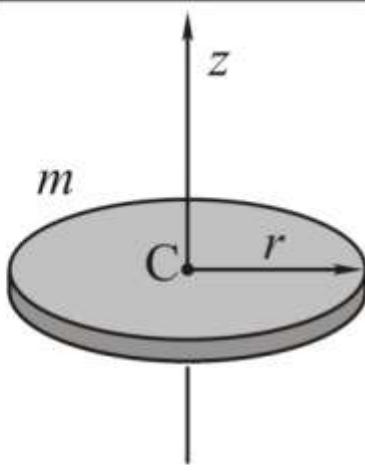
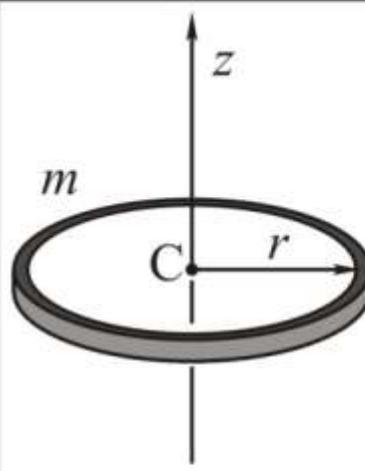
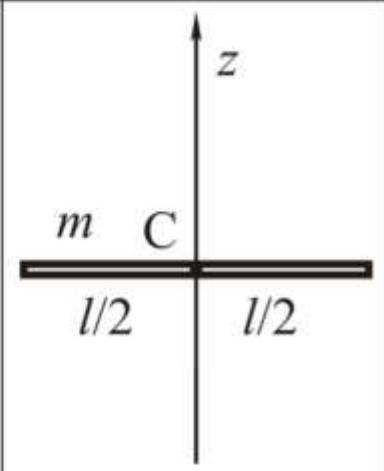
$$\boxed{J_z \ddot{\phi} = \sum M_z^s}$$

- Kinetička energija

$$\boxed{E_K = \frac{1}{2} J_z \omega^2}$$

# Obrtanje oko nepokretne ose

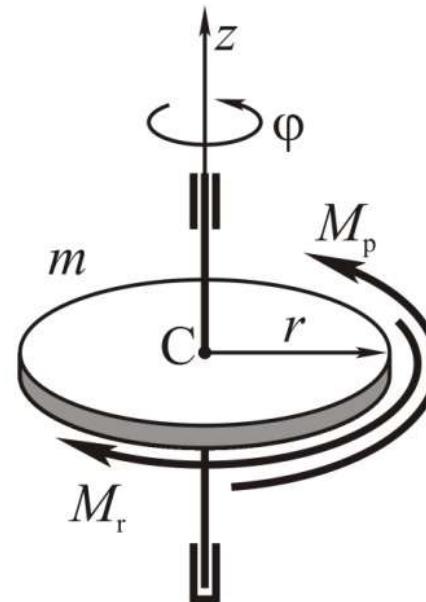
- Moment inercije - primeri

 <p>A diagram of a cylinder of mass <math>m</math> rotating about its central axis <math>z</math>. The center of mass is labeled <math>C</math>, and the radius of rotation is <math>r</math>.</p>	 <p>A diagram of a disk of mass <math>m</math> rotating about its central axis <math>z</math>. The center of mass is labeled <math>C</math>, and the radius of rotation is <math>r</math>.</p>	 <p>A diagram of an elliptical disk of mass <math>m</math> rotating about its central axis <math>z</math>. The center of mass is labeled <math>C</math>, and the radius of rotation is <math>r</math>.</p>	 <p>A diagram of a rectangular plate of mass <math>m</math> rotating about its central axis <math>z</math>. The center of mass is labeled <math>C</math>, and the width of the plate is <math>l</math>.</p>
$J_z = \frac{1}{2}mr^2$	$J_z = \frac{1}{2}mr^2$	$J_z = mr^2$	$J_z = \frac{1}{12}ml^2$

# Obrtanje oko nepokretne ose

- **Primer**

Disk, mase  $m$  i poluprečnika  $r$ , obrće se oko nepokretne ose. Na njega deluje pogonski moment koji je funkcija ugaone brzine –  $M_p = M_0 \left(1 - \frac{\dot{\phi}}{\Omega}\right)$  i konstantan radni moment  $M_r$ . Odrediti zakon promene ugaone brzine diska, ako onkreće iz stanja mirovanja.



$$J_z \ddot{\phi} = M_p - M_r; \quad J_z = \frac{mr^2}{2}$$

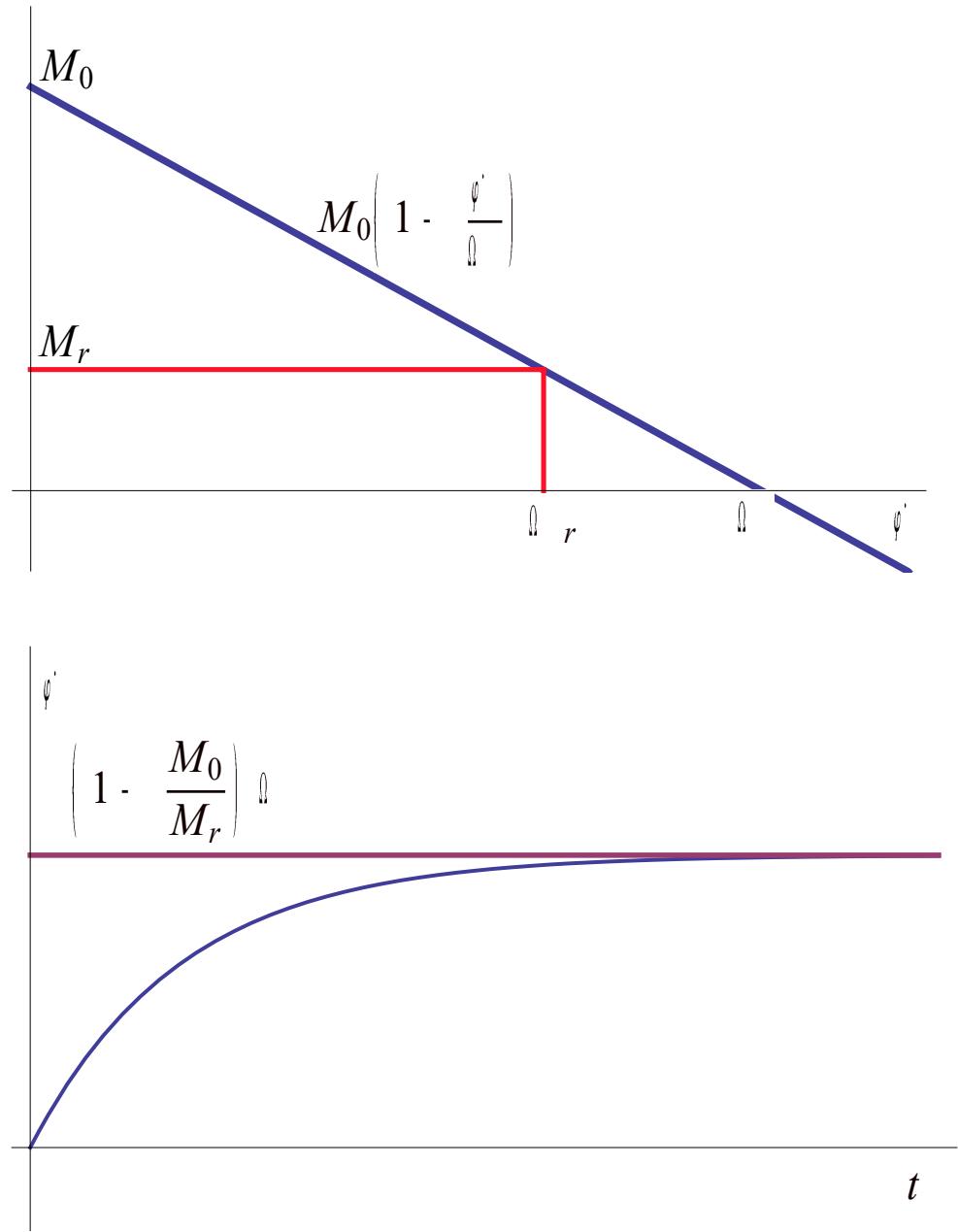
$$J_z \frac{d\dot{\phi}}{dt} = M_0 \left( 1 - \frac{\dot{\phi}}{\Omega} \right) - M_r$$

$$\int_{\dot{\phi}(0)=0}^{\dot{\phi}} \frac{J_z}{M_0 \left( 1 - \frac{\dot{\phi}}{\Omega} \right) - M_r} d\dot{\phi} = \int_0^t dt$$

$$-\frac{J_z \Omega}{M_0} \ln \left( \frac{M_0 \left( 1 - \frac{\dot{\phi}}{\Omega} \right) - M_r}{M_0 - M_r} \right) = t$$

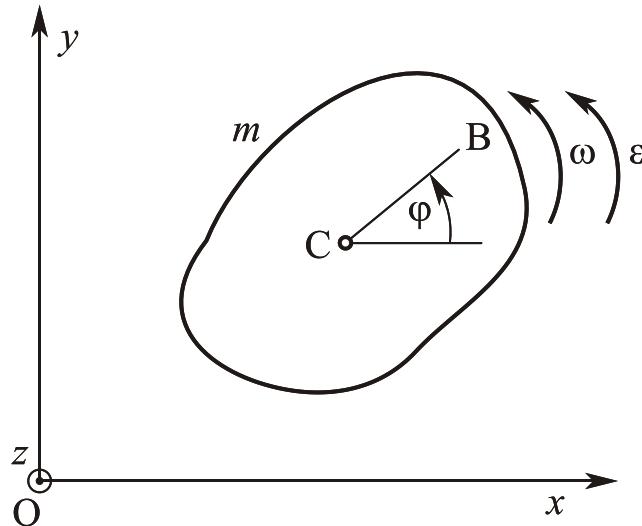
$$\frac{M_0 \left( 1 - \frac{\dot{\phi}}{\Omega} \right) - M_r}{M_0 - M_r} = e^{-\frac{M_0}{J_z \Omega} t}$$

$$\boxed{\dot{\phi} = \Omega \left( 1 - \frac{M_r}{M_0} \right) - \Omega \left( 1 - \frac{M_r}{M_0} \right) e^{-\frac{M_0}{J_z \Omega} t}}$$



# Ravansko kretanje

- Diferencijalne jednačina kretanja



$$m\vec{a}_C = \vec{F}_g^s \rightarrow \begin{cases} m\ddot{x}_C = \sum F_{ix}^s & (1) \\ m\ddot{y}_C = \sum F_{iy}^s & (2) \end{cases}$$

$$\frac{d\vec{L}_C}{dt} = \vec{M}_C^s \rightarrow J_C \ddot{\phi} = \sum M_C^s \quad (3)$$

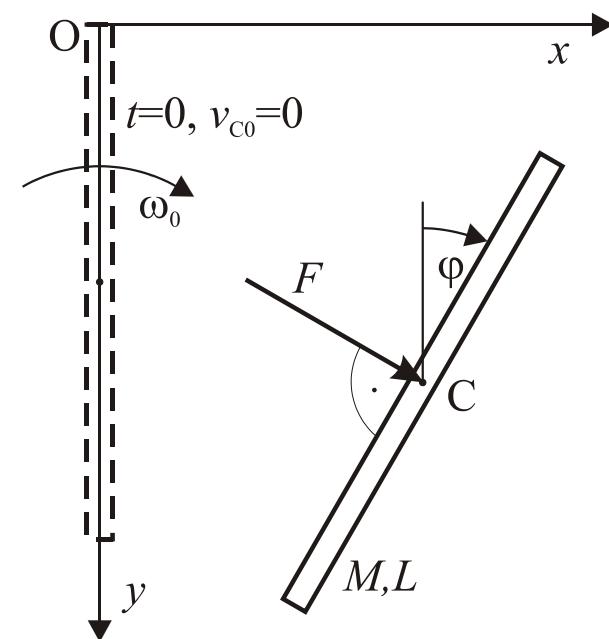
Kinetička energija

$$E_K = \frac{1}{2}mv_C^2 + \frac{1}{2}J_C\omega^2$$

# Ravansko kretanje

- Primer

Štap, mase  $M$  i dužine  $L$ , kreće se po horizontalnom glatkom stolu, pri čemu na njega deluje sila konstantnog intenziteta  $F$ , čija napadna linija stalno prolazi kroz centar štapa i upravna je na njega. Odrediti kretanje štapa,  $(x_C(t), y_C(t), \varphi(t))$ , ako se on u početnom trenutku nalazio na  $y$  osi i ako mu je brzina centra bila jednaka nuli, a ugaona brzina  $\omega_0$ .

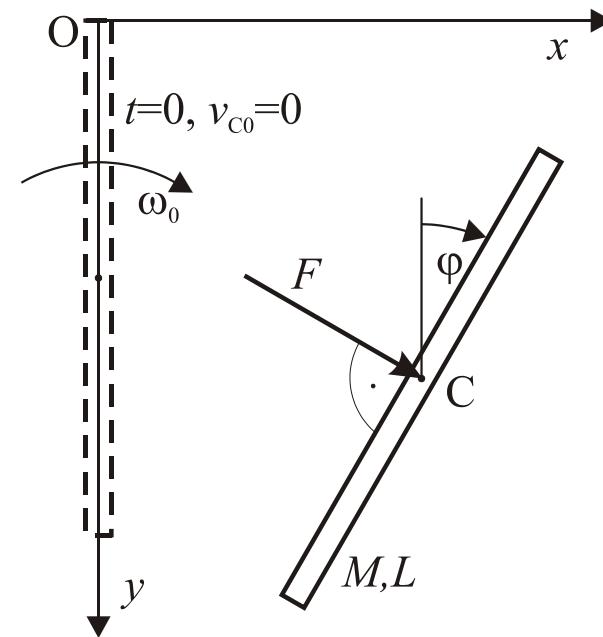


# Ravansko kretanje

- Primer

$$M\vec{a}_C = \vec{F}_g^s \rightarrow \begin{cases} M\ddot{x}_C = F \cos \varphi & (1) \\ M\ddot{y}_C = F \sin \varphi & (2) \end{cases}$$

$$J_C \ddot{\varphi} = \sum M_C^s \rightarrow J_C \ddot{\varphi} = 0 \quad (3)$$



# Ravansko kretanje

- Primer

$$M\vec{a}_C = \vec{F}_g^s \rightarrow \begin{cases} M\ddot{x}_C = F \cos \varphi & (1) \\ M\ddot{y}_C = F \sin \varphi & (2) \end{cases}$$

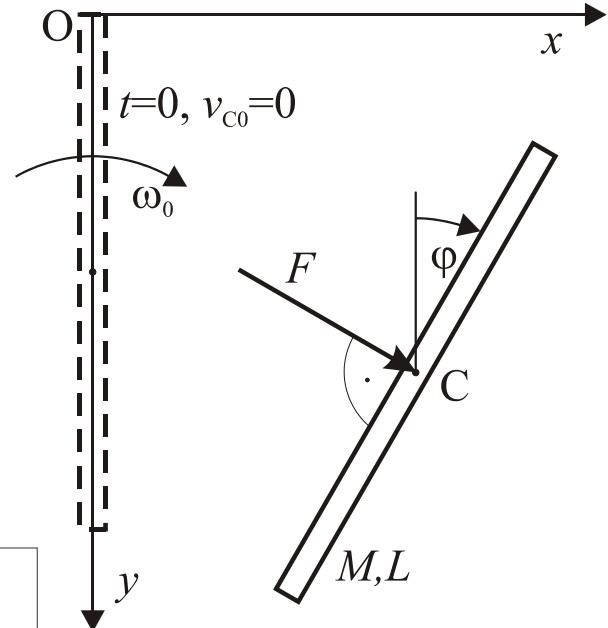
$$J_C \ddot{\varphi} = \sum M_C^s \rightarrow J_C \ddot{\varphi} = 0 \quad (3)$$

$$(3) \rightarrow \ddot{\varphi} = 0 \rightarrow \dot{\varphi} = C = \omega_0 \rightarrow \varphi = \omega_0 t$$

$$(1) \rightarrow M\ddot{x}_C = F \cos \omega_0 t \rightarrow$$

$$\int_0^{x_C} d\dot{x}_C = \frac{F}{M} \int_0^t \cos \omega_0 t \, dt \rightarrow \dot{x}_C = \frac{F}{M\omega_0} \sin \omega_0 t \rightarrow$$

$$\int_0^{x_C} dx_C = \frac{F}{M\omega_0} \int_0^t \sin \omega_0 t \, dt \rightarrow x_C = -\frac{F}{M\omega_0^2} (\cos \omega_0 t - 1)$$



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